Limitations on Precision Cosmology using Clusters of Galaxies

Eric J. Hallman(University of Colorado, Boulder), Patrick M. Motl (Louisiana State University), Jack O. Burns (University of Colorado, Boulder), Michael L. Norman(University of California, San Diego)

Abstrac

We critically analyze the role of clusters of galaxies as probes for precision cosmology. Using synthetic observations of numerically simulated clusters viewed through their X-ray emission and thermal Sunyaev-Zeldovich effect (SZE), we extract measurements of the cluster gas mass. We quantify the possible sources of uncertainty and systematic bias associated with the common simplifying assumptions used in reducing real cluster observations including isothermality and hydrostatic equilibrium. We find that intrinsic variations in clusters limit the precision of observational gas mass estimation to ~10% to 80% confidence excluding instrumental effects. For the full cluster sample, methods that use SZE profiles out to roughly the virial radius are the meat accurate and precise way to estimate cluster mass. X-ray methods are systematically more precise mass estimatos than are SZE methods if merger systems are removed, but X-ray methods slightly overest instance (5-10%), the cluster mass so in aver ago. Finally, we find that methods using a universal temperature profile estimate cluster masses to higher precision than three assuming sothermality.

Introduction

The number density of galaxy clusters as a function of mass and redshift can be used to constrain a number of cosmological parameters. In particular, counting the abundance of clusters above some lower mass limit as a function of cluster redshift places coextraints on Ω_b , Ω_{s_b} , Ω_{s_b} , and the dark energy equation of state, w (Pen (1997), Sasaki (1996)). Measuring the apparent change of cluster gas fraction with redshift also produces constraints on w (e.g. Vikh limit et al. (2005)). In order for such methods to provide strong constraints on cosmological parameters, cluster mass estimators must be accurate to \sim 10% (Haiman et al. 2001). Previous studies (e.g. Evrand et al. (1996)) suggests that this level of precision may be difficult to achieve. The often used scaling relation between cluster mass and X-ray spectral temperature has relatively small scatter compared to other methods, but still is uncertain to \sim 30% (e.g. Sanderson et al. (2003)). This study aims to determinate the limiting accuracy of various X-ray and Sunyaev-Zeldovich effect (SZE) observational methods of cluster gas mass estimation. This determination will indicate whether cluster methods can be accurate enough to do precision comology. This study addresses the key question: What is the best way to massure cluster gas masses with precision in order to do associated.

High resolution X-ray or SZE observations of clusters coupled with assumptions about the gas distribution lead to estimates of the gas mass in the cluster dark matter potential well. The electron number density is often assumed to fit an isothermal β model,

$$n_c(r) = n_{c0} \left(1 + \left(\frac{r}{r_c}\right)^2\right)^{-3\beta/2}$$
. (1)

Fitting an observed X-ray or SZE profile to these projected isothermal β model X-ray surface brightness and y parameter distributions results in a description of the density distribution, which can be integrated to obtain the gas mass. The difference in dependence on gas density and temperature of X-ray emissivity and the SZE y parameter makes the combination of these two methods of observing tools results in the size of observing clusters not only observations of individual clusters allow one to extract the density and temperature of the gss without relying on X-ray spectral temperatures.

Recent numerical (e.g. Loken et al. (2002)) and observational studies (e.g. Vikhlinin et al. (2005)) suggest that the cluster ggs is not isothermal, but fits more closely to a universal temperature profile (UTP), for this study written as

$$T(r) = \langle T \rangle_{500} T_0 \left(1 + \left(\frac{r}{\alpha r_{500}} \right)^2 \right)^{-\delta},$$
 (

where $(T)_{500}$ is the average temperature inside a projected overdensity radius of r_{500} . T_{0} , α , and δ are parameters whose mean value is measured from the entire cluster population at each redshift. We perform a geometric deprojection of the X-ray and SZE profiles and use this additional assumption about the gas temperature to calculate the density profile and the mass.

With observations of both X-ray and SZE for a particular cluster, an observer can deproject the surface brightness of each simultaneously to determine the gas density. This method is particularly powerful due to the difference in dependence of X-ray and SZE ensists on a density and temperature. In this study, we also perform such a combined deprojection of X-ray and SZE profiles to obtain the gas density radial profile with no assumption about the temperature profile

We examine the effect of the assumption of the gas temperature and density profile on the determination of the true cluster gas mass. Using adaptive melh hydro/N-body cosmological simulations, we have extracted all clusters with $M \geq 10^{14} M_{\odot}$ out to a redshift of y = 2. The four simulations used progressively more explaint baryonic physics. For each cluster, we have fit the X-ray and SZEs urince brightness profiles to those produced by a base-lift isothermal gas distribution. We also perform a direct geometric deprojection of the SZE and X-ray

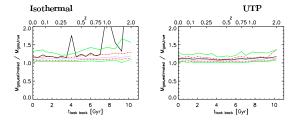
Our simulations use the hybrid Eulerian adaptive mesh refinement/n-body code Enzo (O'She et al. (2005); http://cosmo-uscd-edn/enzo) to evolve both the dark matter and baryonic fluid in the clusters; utilizing the piecewise parabolic scheme for the hydrodynamics. With up to seven levels of dynamic refinement in high density regions, we attain spatial resolution up to ~ 16 kpc in the clusters. We assume a cotoordance ~ 100 Mccsmological model with the following parameters: $\Omega_b = 0.026$, $\Omega_b = 0.3$, $\Omega_b = 0.7$, h = 0.7, and $\sigma_b = 0.928$.

Table 1. Ratio of Estimated Gas Mass to True Gas Mass.

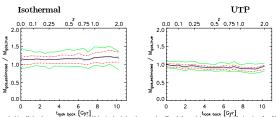
100	Table 1. Table of Edifficated Gas Wash to Trac Gas Wash					TYRONDEA
Method	Model	Radius	Median	Mean	80% Upper	80% Lower
Xray	UTP	r_{500}	1.06	1.07	1.11	1.03
Xray	Iso	r_{500}	1.08	1.08	1.13	1.04
Joint	Geometric	r_{500}	1.09	1.09	1.16	1.02
SZE	UTP	r_{200}	0.98	0.99	1.06	0.93
SZE	UTP	r_{500}	1.01	1.01	1.11	0.92
SZE	Iso	r_{500}	1.08	1.07	1.20	0.96
SZE	y-M	r_{500}	0.96	0.98	1.13	0.87
Xray	T_X - \mathbf{M}	r_{500}	1.03	1.01	1.27	0.77

Comparison of Methods as a Function of Redshift

X-Ray



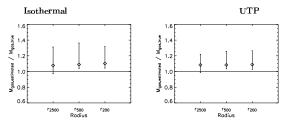
SZE



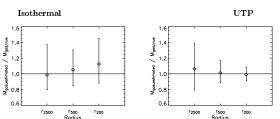
To be so to

Comparison of Estimates at Different Radii

X-Ray



SZE



Hence 2. Top: Left panel shows modals ratio of estimated gas mass to two gas mass inside various and as indicated for all clusters above 10.1 Mg, at x=0 in the simulation for X-ray insthemal mass estimate. Error bars are SSW scatter. Ugin panel is similar plot for deprojection of X-ray assuring a UTF modal for temperature Bottom: Left panel is mass ratio for friging SZD y parameter destination to implement model, edgis panel is similar but for SZ deposition assuring a UTF compensation. Hence shows are present Mag. = Mag.

X-ray vs. SZE Mass Estimation

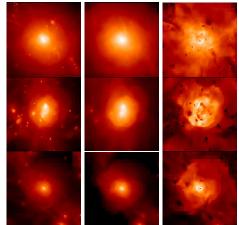


Figure 3. Three clusters. Top row appears relaxed, middle row has disturbed morphology and bottom row has a cool core. Left column shows X-ray surface brightness from simulated cluster, middle column is Compton y parameter and right column is emission weighted temperature.

Table 2. Gas Mass Estimates for Three Clusters

Cluster State	X-ray M_{est}/M_{true}	SZE M_{cst}/M_{true}	Total Mass (M_{\odot})
Relaxed	1.00	0.94	1.5×10^{15}
Dist ur bed	1.54	1.07	5.3×10^{14}
Cooling Core	1.39	1.10	1.4×10^{14}

Summary

The highest precision to which cluster gas masses can be measured using typical assumptions and either X-ray, SZE or a combination of both measurements is $\sim 10\%$ to 80% confidence. This study does not include instrumental or other observational effects, and so is an indication of the limiting ability of observations to correctly gauge the mass of clusters in the ideal case. These limits in precision are a direct result of the deviation of the simulated clusters from simple assumptions about their physical and thermodynamic properties, dynamical state, and sphericity.

Mass estimates fitting radial profiles to a universal temperature profile (UTP) have smaller scatter than similar estimates assuming isothermality.

SZE methods of gos mass estimation assuming a universal temperature profile in the cluster gos produce the smallest scatter when estimating masses in a raw sample of clusters. Cleaning the cluster sample for disturbed or merging clusters is much less important in SZE methods, particularly when profiles can be observed to large radius. As a practical matter, SZE methods are superior for mass estimation for large samples of clusters out to high redshift. The SZE is more representative of the true-cluster potential than the X-ray emission.

X-ray methods show a 5-10% bias in median values which is absent is SZE methods. The bias is a result of substructure and mergers in clusters, events which enhance the gas density locally. The stronger density dependence of X-ray emissivity tends to drive up the X-ray mess estimates more than the SZE estimates.

Cool core clusters in our catalogs are particularly poor candidates for precision mass estimation. Even when excising the cool core from the analysis, mass estimation shows larger scatter in both X-ray and SZE methods than for non-cool core clusters. The emission profiles of X-ray and SZE are affected outside the cool core rescion due to the presence of cold substructure in the clusters.

References

Evrard, A. E., Metzler, C. A., & Navarro, J. F. 1996, ApJ, 469, 494.

Haiman, Z., Mohr, J. J., & Holder, G. P. 2001, ApJ, 553, 545.

Loken, C., Norman, M. L., Nelson, E., Burns, J., Bryan, G. L., & Motl, P. 2002, ApJ, 579, 571.
O'Shea, B. W., et al. 2005. in Adaptive Mesh Refinement: Theory and Applications. Berlin: Springer

Pen, U. 1997, New Astronomy, 2, 309

Sanderson, A. J. R., Pomnan, T. J., Finoguenov, A., Lloyd-Davies, E. J., & Markevitch, M. 2003, MNRAS, 340, 989. Sasaki, S. 1996, PASJ, 48, L119.

Sasaki, S. 1996, PASJ, 48, I.119.
Vikhlinin, A., Markeyitch, M., Murray, S. S., Jones, C., Fonnan, W., & Van Speybroock, L. 2005, ApJ, 628, 655.

Acknowledgments

The simulations presented in this paper were conducted at the National Center for Supercomputing Applications at the University of Illinois, Urbana-Champaign. We acknowledge the support of the Chandra X-ray Center and NASA through grant TM3-0938. We wish to acknowledge the support of the National Science Foundation through grant AST-0407368. We also wish to thank the Laboratory for Computational Astrophysics for support of the Research and the Computational Astrophysics for support of the Research and the Computational Astrophysics for support of the Research and the Computational Astrophysics for support of the Research and the Computational Astrophysics for support of the Research and the Computational Astrophysics for support of the Research and the Computational Astrophysics for support of the Research and the Computational Astrophysics for support of the Research and Rese

This miniposter was prepared with Brian Wolven's Poster LATEX macros v2.1.

Table 1:

Table 2:

Fig. 1.—

Fig. 2.—

Fig. 3.—

Fig. 4.—

ъ. -