Rapidly rotating black holes in Chern-Simons gravity: Decoupling limit solutions and breakdown

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Einstein Fellows Symposium, Cambridge, MA, 29 Oct. 2014

Phys. Rev. D 90, 044061 (2014) [arXiv:1407.2350]

Motivation

1 GR successful but incomplete

- GR+QM=new physics (e.g. BH thermo)
- Planck scale phenomena? Other scales?
- Expect GR is low-energy EFT
- Ø Ask nature
 - So far, only weak-field tests
 - Lots of theories $\approx {\rm GR}$
 - Need to explore strong-field
 - Strong curvature non-linear





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What phenomena come from UV completions?

Theories

Fundamental approach:

- String theory Loop quantum gravity
- TeVeS Einstein-Æther Hŏrava
- Massive gravity dRGT bi-metric

• . . .

Pedestrian approach: effective field theory

- Learned from cond-mat, then nuclear and hep-th
- Theory with separation of scales
- "Integrate out," effective theory for long (or short) wavelengths
- Works backwards!



Theories

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Worked for describing superconductivity, predicting W, higgs, etc.

Try to build EFT for gravity

- Metric, general covariance, Lorentz invariance
- Lowest order dynamical theory is $\Lambda + {\sf GR}!$

Beyond GR: add new ℓ —want to constrain this

Dynamical Chern-Simons Gravity

$$S = \int d^4x \sqrt{-g} \left[\frac{m_{\rm pl}^2}{2} R - \frac{1}{2} (\partial \vartheta)^2 + \frac{m_{\rm pl}}{8} \ell^2 \vartheta \ ^*\!RR \right]$$

- Anomaly cancellation, low-E string theory, ...
- Lowest-order EFT with parity-odd ϑ , shift symmetry (long range)
- Phenomenology unique from other R² (e.g. Einstein-dilaton-Gauss-Bonnet)
- Tractable
- Straw-man theory

- Most EFTs don't make sense as exact theories (see e.g. Delsate+Hilditch+Witek)
- (Almost) All corrections introduce new ℓ
- Can't be too long
- Expand fields, EOMs in powers of $\ell/\mathcal{R}_{\rm BG}$, perturbation scheme

Question: What is regime of validity of decoupling limit?

How do we constrain ℓ in dCS from astronomical observations?

• dCS is higher-curvature and parity-odd



How do we constrain ℓ in dCS from astronomical observations?

dCS is higher-curvature and parity-odd

$$\Box \vartheta = -\frac{m_{\rm pl}}{8} \ell^2 \ ^*\!RR \,,$$

- Want *RR as large as possible \implies smallest M, largest $\chi = J/M^2$
- NSs have small M, but BHs have $\chi \to 1$

Want solutions for rapidly rotating black holes in dCS

- a = 0 (Schwarzschild) is exact solution with $\vartheta = 0$
- Analytically known solutions in decoupling limit
 - $a \ll M$ limit up to $\mathcal{O}(a^2)$, valid $\forall r$ (see Yunes+Pretorius, Yagi+Yunes+Tanaka)
 - $r \gg M$ limit for l = 1, valid $\forall a$ (see Yagi+Yunes+Tanaka)
- I construct numerical solutions $\forall r, \forall a$

Takeaway

- I construct numerical solutions $\forall r, \forall a \text{ for dCS BHs in decoupling limit}$
- I use solutions to determine the regime of validity of PT



- This can be turned around to forecast bounds $\ell \lesssim 22$ km from GRO J1655–40 ($M = 6.30 \pm 0.27 M_{\odot}$, $\tilde{a} \approx 0.65$ –0.75)
- For details see Phys. Rev. D 90, 044061 (2014) [arXiv:1407.2350]

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Equations to solve

$$S = \int d^4x \sqrt{-g} \left[\frac{m_{\rm pl}^2}{2} R - \frac{1}{2} (\partial \vartheta)^2 + \frac{m_{\rm pl}}{8} \ell^2 \vartheta \ ^*RR \right]$$
$$\Box \vartheta = -\frac{m_{\rm pl}}{8} \ell^2 \ ^*RR$$
$$m_{\rm pl}^2 G_{ab} + m_{\rm pl} \ell^2 C_{ab} = T_{ab}^{(m)} + T_{ab}^{(\vartheta)}$$

•
$$g^{ab}C_{ab} = 0$$

Take $\ell^2 \rightarrow \varepsilon \ell^2$ and expand in ε ,

$$artheta = 0 + arepsilon artheta^{(1)} + rac{arepsilon^2}{2}artheta^{(2)} + \dots$$

 $g = g_{
m GR} + arepsilon 0 + arepsilon^2 h^{
m def} + \dots$
 $\Box_{
m GR} artheta^{(1)} = -rac{m_{
m pl}}{8} \ell^2 \ ^*\!RR[g_{
m GR}]$

Don't yet have quantity to test validity of perturbation theory

Next order

$$m_{\rm pl}^2 G_{ab}^{(1)}[h^{\rm def}] = T_{ab}^{(\vartheta)}[\vartheta^{(1)}, \vartheta^{(1)}] - m_{\rm pl} \ell^2 C_{ab}[\vartheta^{(1)}]$$

- Trace equation
- Lorenz gauge $\nabla_a \bar{h}^{ab} = 0$

$$\frac{1}{2}m_{\rm pl}^2\Box h^{\rm def} = -(\nabla^a\vartheta^{(1)})(\nabla_a\vartheta^{(1)})\,,$$

- Same scalar PDE operator
- Caveat: gauge-dependent but should still capture *a*-dependence

• Now can make comparison:

$$\sqrt{-g} = \sqrt{-g_{\rm GR}} \left(1 + \varepsilon^2 \frac{1}{2} h^{\rm def} + \mathcal{O}(\varepsilon^3) \right)$$

- If $h^{\mathrm{def}} \sim \mathcal{O}(1)$, should keep higher $\mathcal{O}(\varepsilon)$
- Criterion for validity of PT:

$$|h^{
m def}| \lesssim 1$$
 everywhere

• Program: Solve for $\vartheta^{(1)}, h^{\text{def}}$ as functions of r, θ for all a

Approach to solving

Symmetry reduced, $\vartheta = \vartheta(r, \theta)$. $\Box \to \Delta$.

Analytical:

- Static Green's function Δ^{-1} known analytically
- Separation of variables

$$\vartheta = \sum_{j} \vartheta_j(r) P_j(\cos \theta)$$

• Can do source decomposition (See Konno+Takahashi and [arXiv:1407.0744])

Resort to numerics!

- Elliptic PDE. Could solve hyperbolic, parabolic, relaxation scheme
- Numerical separation of variables. Each j mode is an ODE.
- Compactify r
- Pseudospectral collocation method
- Directly solve discrete ODE operator ("numerical Green's function")

Numerical approach

- For each a, find $\vartheta(r,\theta;a)$, compute $(\partial \vartheta)^2$, find $h^{\text{def}}(r,\theta;a)$
- Evaluate $\max |h^{ ext{def}}|$ and find regime of validity





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Exponential convergence



Exponential convergence



Exponential convergence



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Rapidly rotating BHs in dCS

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 $\tilde{\mathbf{r}}\sin\theta$

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Regime of validity



Forecasting bounds

- Observation of BH indistinguishable from GR predictions
- Size of ℓ correction below breakdown (caveat: cancellation)
- GRO J1655–40: $M = 6.30 \pm 0.27 M_{\odot}$, $\tilde{a} \approx 0.65$ –0.75



• Better by 10^7 than Solar System bounds

- $a \to GM$ limit analytically?
- All a analytically?
- Rest of the metric
- Accretion disk modeling

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