

# Rapidly rotating black holes in Chern-Simons gravity: Decoupling limit solutions and breakdown

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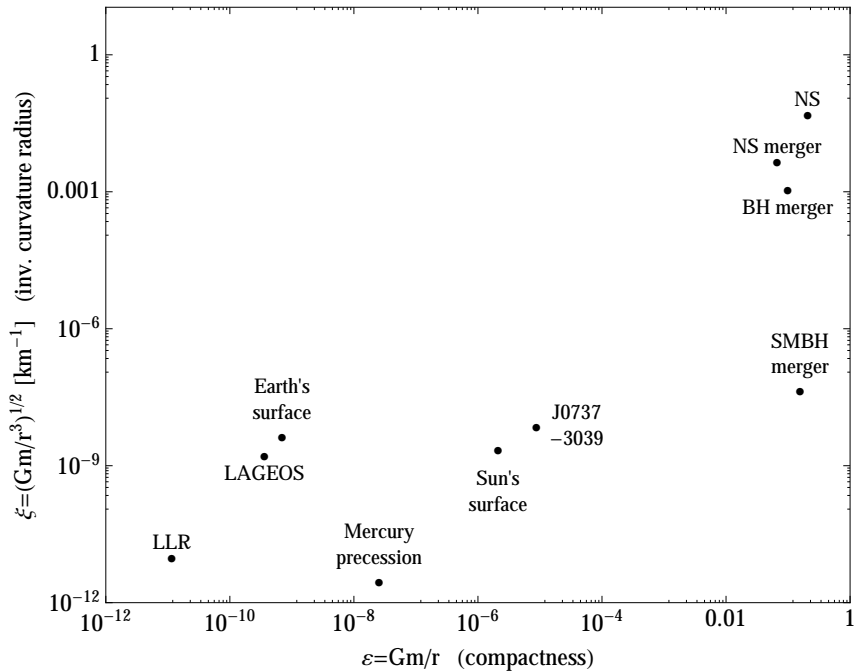
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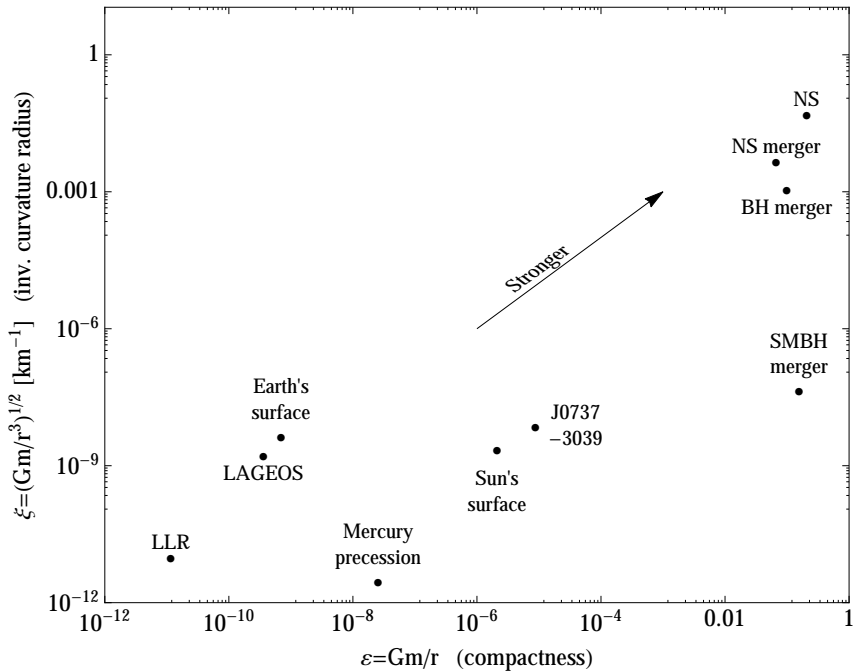
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Phys. Rev. D 90, 044061 (2014) [arXiv:1407.2350]

# Motivation

- ① GR successful but incomplete
  - GR+QM=new physics (e.g. BH thermo)
  - Planck scale phenomena? Other scales?
  - Expect GR is low-energy EFT
- ② Ask nature
  - So far, only weak-field tests
  - Lots of theories  $\approx$  GR
  - Need to explore strong-field
    - Strong curvature • non-linear





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What phenomena come from UV completions?

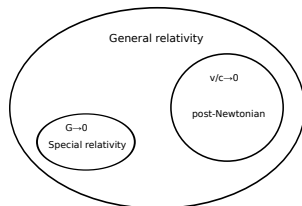
# Theories

Fundamental approach:

- String theory • Loop quantum gravity
- TeVeS • Einstein-Æther • Hřrava
- Massive gravity • dRGT • bi-metric
- ...

Pedestrian approach: effective field theory

- Learned from cond-mat, then nuclear and hep-th
- Theory with separation of scales
- “Integrate out,” effective theory for long (or short) wavelengths
- Works backwards!



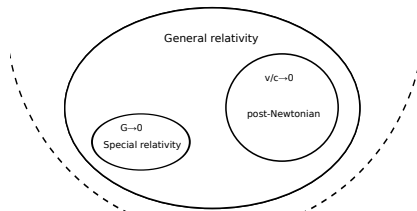
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# EFT works

Worked for describing superconductivity, predicting W, higgs, etc.

Try to build EFT for gravity

- Metric, general covariance, Lorentz invariance
- Lowest order dynamical theory is  $\Lambda$ +GR!

Beyond GR: add new  $\ell$ —want to constrain this



# Dynamical Chern-Simons Gravity

$$S = \int d^4x \sqrt{-g} \left[ \frac{m_{\text{Pl}}^2}{2} R - \frac{1}{2} (\partial\vartheta)^2 + \frac{m_{\text{Pl}}}{8} \ell^2 \vartheta *RR \right]$$

- Anomaly cancellation, low-E string theory, ...
- Lowest-order EFT with parity-odd  $\vartheta$ , shift symmetry (long range)
- Phenomenology unique from other  $R^2$  (e.g. Einstein-dilaton-Gauss-Bonnet)
- Tractable
- Straw-man theory

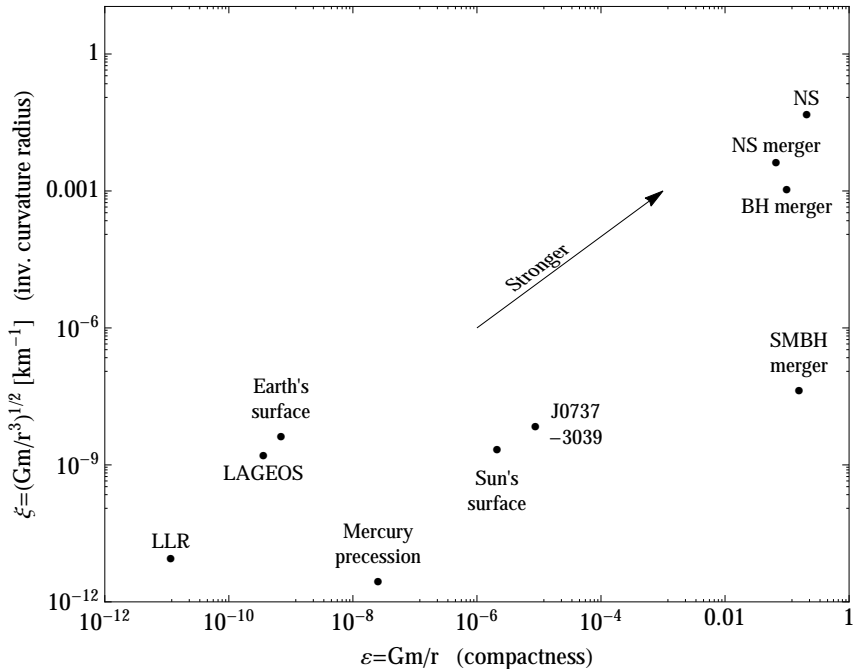
# Decoupling limit

- Most EFTs don't make sense as exact theories (see e.g. Delsate+Hilditch+Witek)
- (Almost) All corrections introduce new  $\ell$
- Can't be too long
- Expand fields, EOMs in powers of  $\ell/\mathcal{R}_{\text{BG}}$ , perturbation scheme

Question: What is regime of validity of decoupling limit?

# How do we constrain $\ell$ in dCS from astronomical observations?

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- dCS is higher-curvature and parity-odd

$$\square\vartheta = -\frac{m_{\text{pl}}}{8}\ell^2 {}^*RR,$$

- Want  ${}^*RR$  as large as possible  
 $\implies$  smallest  $M$ , largest  $\chi = J/M^2$
- NSs have small  $M$ , but BHs have  $\chi \rightarrow 1$

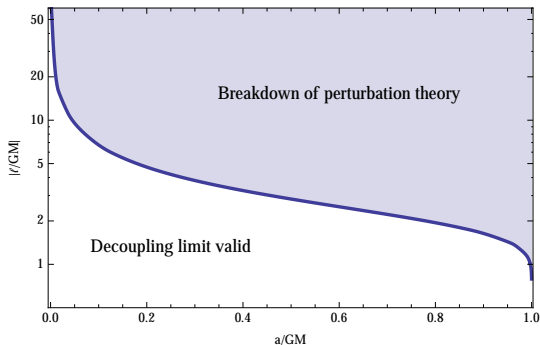
Want solutions for rapidly rotating black holes in dCS

# Black holes in dCS

- $a = 0$  (Schwarzschild) is exact solution with  $\vartheta = 0$
- Analytically known solutions in decoupling limit
  - $a \ll M$  limit up to  $\mathcal{O}(a^2)$ , valid  $\forall r$  (see Yunes+Pretorius, Yagi+Yunes+Tanaka)
  - $r \gg M$  limit for  $l = 1$ , valid  $\forall a$  (see Yagi+Yunes+Tanaka)
- I construct numerical solutions  $\forall r, \forall a$

# Takeaway

- I construct numerical solutions  $\forall r, \forall a$  for dCS BHs in decoupling limit
- I use solutions to determine the regime of validity of PT



- This can be turned around to forecast bounds  $\ell \lesssim 22\text{km}$  from GRO J1655–40 ( $M = 6.30 \pm 0.27M_{\odot}$ ,  $\tilde{a} \approx 0.65\text{--}0.75$ )
- For details see [Phys. Rev. D 90, 044061 \(2014\) \[arXiv:1407.2350\]](#)

# Equations to solve

$$S = \int d^4x \sqrt{-g} \left[ \frac{m_{\text{pl}}^2}{2} R - \frac{1}{2} (\partial\vartheta)^2 + \frac{m_{\text{pl}}}{8} \ell^2 \vartheta {}^*RR \right]$$

$$\square\vartheta = -\frac{m_{\text{pl}}}{8} \ell^2 {}^*RR$$

$$m_{\text{pl}}^2 G_{ab} + m_{\text{pl}} \ell^2 C_{ab} = T_{ab}^{(m)} + T_{ab}^{(\vartheta)}$$

- $g^{ab} C_{ab} = 0$



## Decoupling limit

Take  $\ell^2 \rightarrow \varepsilon \ell^2$  and expand in  $\varepsilon$ ,

$$\vartheta = 0 + \varepsilon \vartheta^{(1)} + \frac{\varepsilon^2}{2} \vartheta^{(2)} + \dots$$

$$g = g_{\text{GR}} + \varepsilon 0 + \varepsilon^2 h^{\text{def}} + \dots$$

$$\square_{\text{GR}} \vartheta^{(1)} = -\frac{m_{\text{pl}}}{8} \ell^2 *RR[g_{\text{GR}}]$$

Don't yet have quantity to test validity of perturbation theory

## Next order

$$m_{\text{pl}}^2 G_{ab}^{(1)}[h^{\text{def}}] = T_{ab}^{(\vartheta)}[\vartheta^{(1)}, \vartheta^{(1)}] - m_{\text{pl}} \ell^2 C_{ab}[\vartheta^{(1)}]$$

- Trace equation
- Lorenz gauge  $\nabla_a \bar{h}^{ab} = 0$

$$\frac{1}{2} m_{\text{pl}}^2 \square h^{\text{def}} = -(\nabla^a \vartheta^{(1)})(\nabla_a \vartheta^{(1)}),$$

- Same scalar PDE operator
- Caveat: gauge-dependent but should still capture  $a$ -dependence

## Next order

- Now can make comparison:

$$\sqrt{-g} = \sqrt{-g_{\text{GR}}} \left(1 + \varepsilon^2 \frac{1}{2} h^{\text{def}} + \mathcal{O}(\varepsilon^3)\right)$$

- If  $h^{\text{def}} \sim \mathcal{O}(1)$ , should keep higher  $\mathcal{O}(\varepsilon)$
- Criterion for validity of PT:

$$|h^{\text{def}}| \lesssim 1 \quad \text{everywhere}$$

- Program: Solve for  $\vartheta^{(1)}, h^{\text{def}}$  as functions of  $r, \theta$  for all  $a$

# Approach to solving

Symmetry reduced,  $\vartheta = \vartheta(r, \theta)$ .  $\square \rightarrow \Delta$ .

Analytical:

- Static Green's function  $\Delta^{-1}$  known analytically
- Separation of variables

$$\vartheta = \sum_j \vartheta_j(r) P_j(\cos \theta)$$

- Can do source decomposition (See Konno+Takahashi and [arXiv:1407.0744])

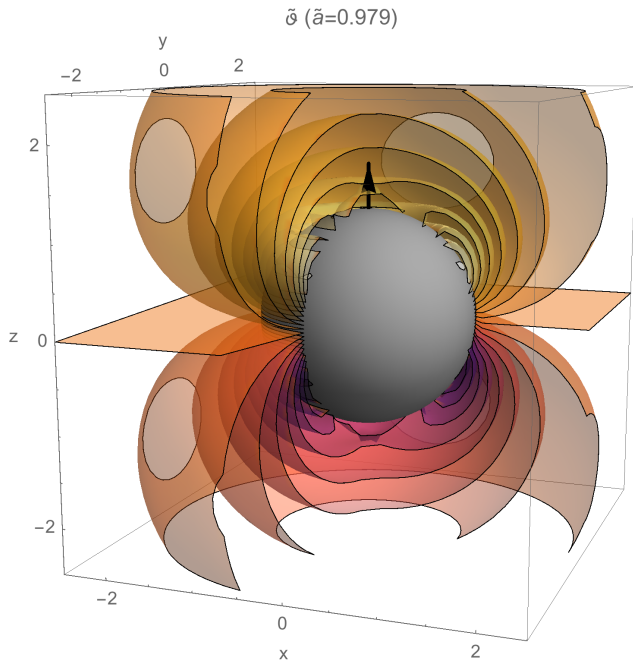
Resort to numerics!

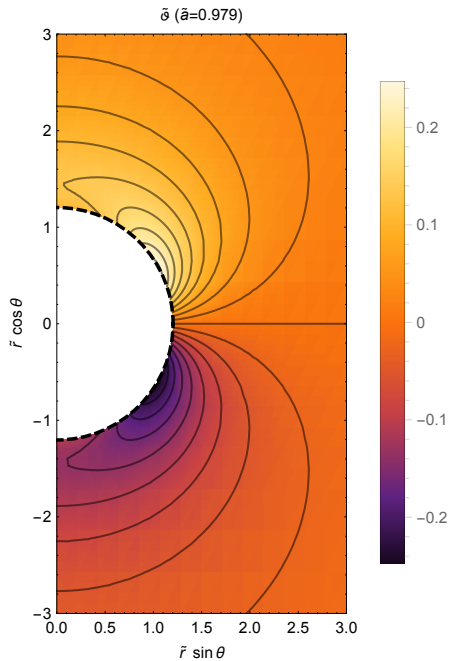
# Numerical approach

- Elliptic PDE. Could solve hyperbolic, parabolic, relaxation scheme
- Numerical separation of variables. Each  $j$  mode is an ODE.
- Compactify  $r$
- Pseudospectral collocation method
- Directly solve discrete ODE operator (“numerical Green’s function”)

# Numerical approach

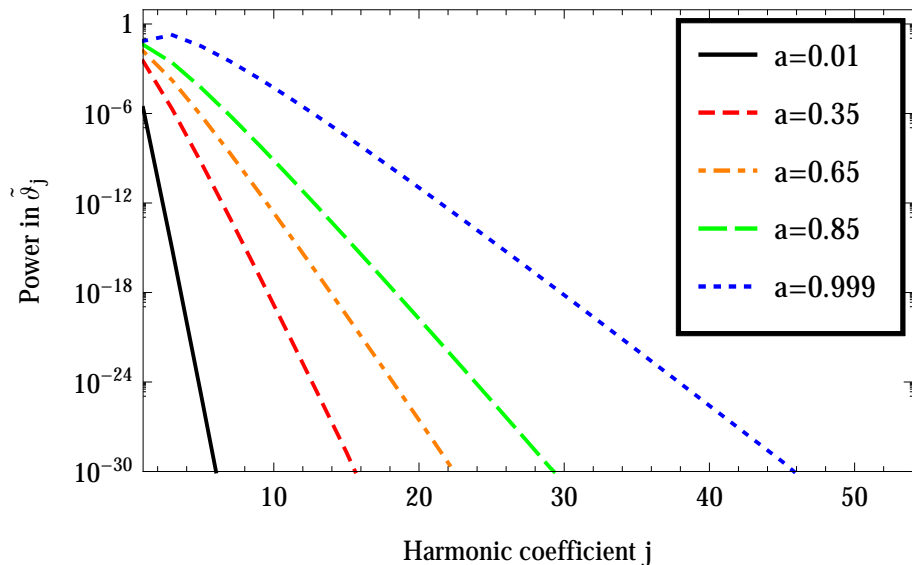
- For each  $a$ , find  $\vartheta(r, \theta; a)$ , compute  $(\partial\vartheta)^2$ , find  $h^{\text{def}}(r, \theta; a)$
- Evaluate  $\max |h^{\text{def}}|$  and find regime of validity



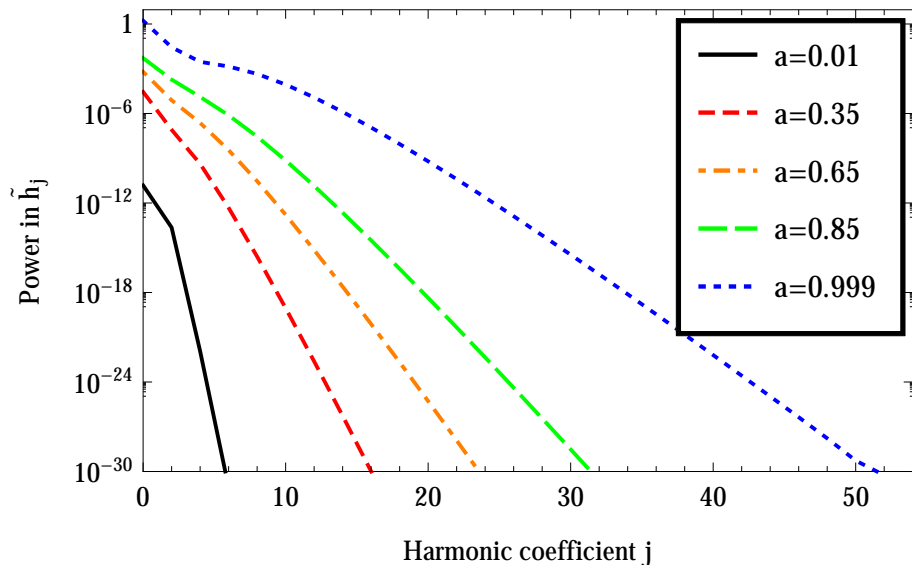




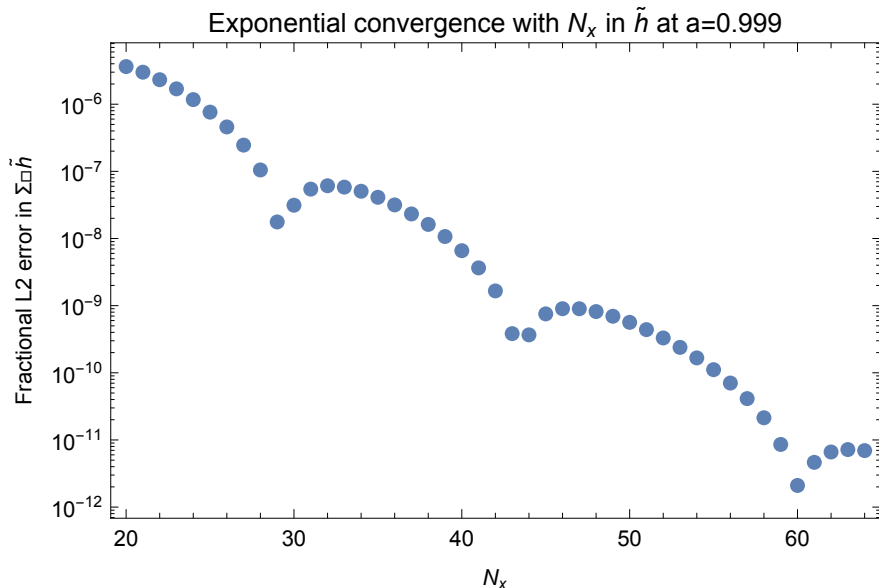
# Exponential convergence

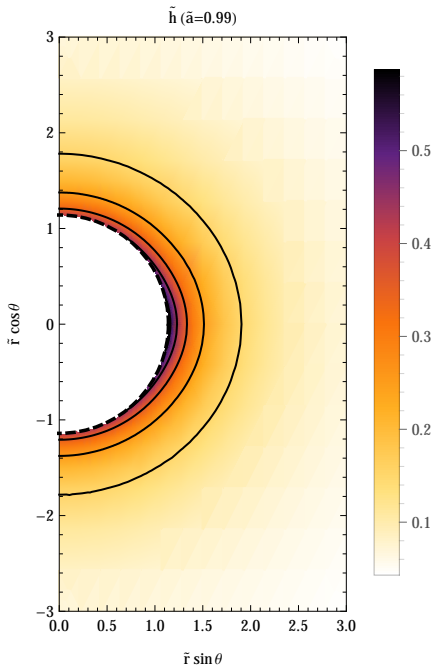


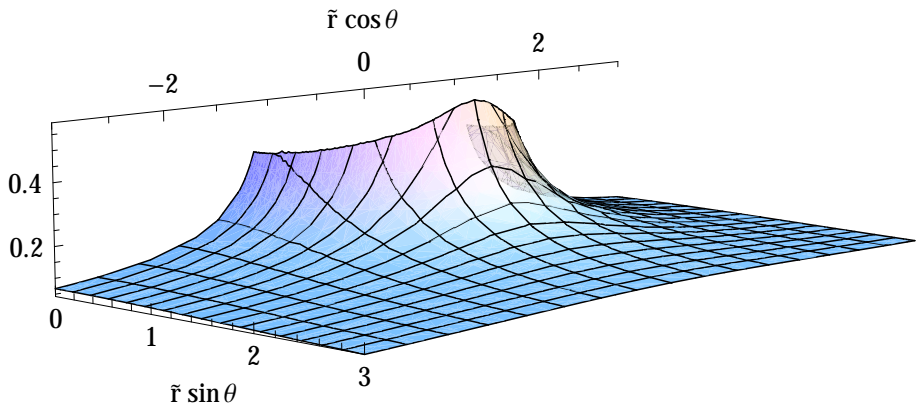
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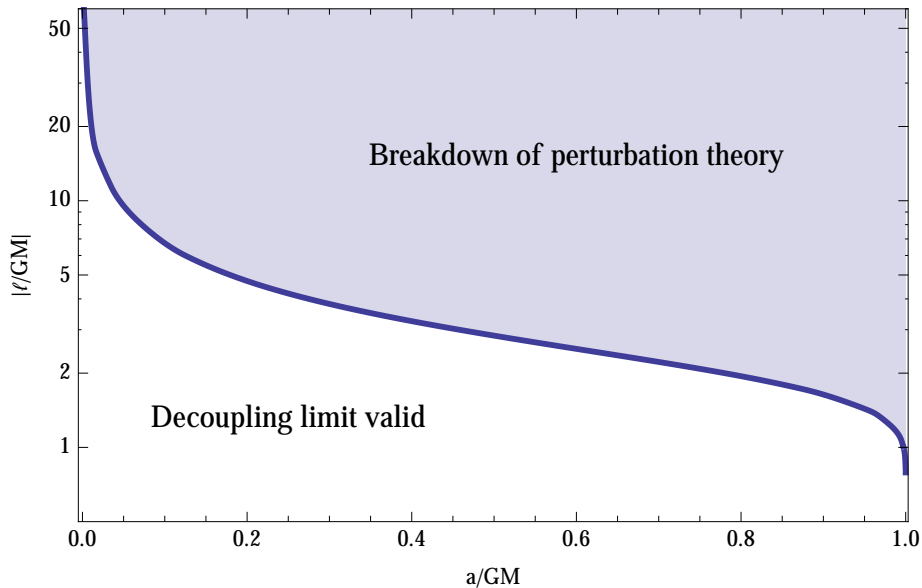
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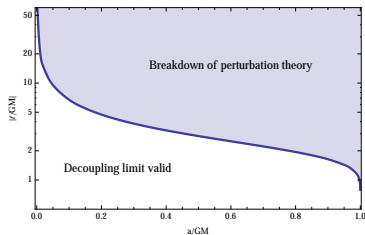


## Regime of validity



# Forecasting bounds

- Observation of BH indistinguishable from GR predictions
- Size of  $\ell$  correction below breakdown (caveat: cancellation)
- GRO J1655–40:  $M = 6.30 \pm 0.27 M_{\odot}$ ,  $\tilde{a} \approx 0.65\text{--}0.75$



$$\Rightarrow \ell \lesssim 22\text{km}$$

- Better by  $10^7$  than Solar System bounds

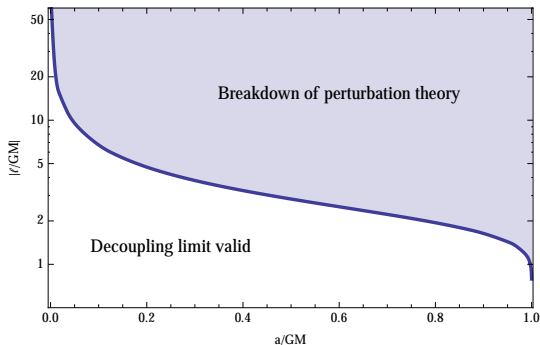
# Future work

- $a \rightarrow GM$  limit analytically?
- All  $a$  analytically?
- Rest of the metric
- Accretion disk modeling



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