# "Two Types of X-ray Spectra in Cataclysmic Variables" Revisited Koji Mukai NASA/GSFC/CRESST and UMBC

**ABSTRACT:** In a 2003 paper, Mukai et al. presented the Chandra HETG spectra of 7 cataclysmic variables (CVs) then available in the public archive, and classified them into "cooling flow" and "photoionized" types. In this presentation, I will revisit this classification. It is clear that multi-temperature plasma exists in the post-shock region of accreting white dwarfs; I will discuss why the cooling flow model works well, and what its limitations are when applied to CVs. It is also clear that even the "photoionized" CVs are also powered by the same multi-temperature plasma. I will discuss how the appearance of a much harder continuum of these "photoionized" CVs is created by the complex intrinsic absorber. I will show that there is an additional soft emission in these systems that the multi-temperature plasma models cannot explain, and discuss what evidence there is for the photoionization origin of this component.



1. Introduction

Cataclysmic variables (CVs) are interacting binaries in which the accreting object (the primary) is a white dwarf (see Warner 1995 for a review). Xray emission in CVs is powered by the accretion of material onto the white dwarf, which is shockheated to high temperatures  $(kT_{\rm max} \sim 10-50 \text{ keV})$ and must cool before settling onto the surface of the primary. The emergent X-ray spectrum is expected to be the sum of emission from plasmas over a continuous temperature distribution, from the shock temperature to the white dwarf photospheric temperature. However, Mukai et al. (2003) presented Chandra HETG spectra of 7 CVs then available in the archive, and showed that they can be divided into two types, "cooling-flow" (Figure 1) and "photoionized" (Figure 2).





### 2. Why Cooling Flow?

The post-shock plasma in CVs loses energy by emitting X-rays and therefore must cool, and cannot remain at a single temperature. In the case of magnetic CVs, one can solve for the temperaturedensity structure of the postshock region assuming pure Bremsstrahlung cooling (Figure 3). The case of non-magnetic CVs accreting via a boundary layer is more complex, but a similar multi-temperature structure is expected. Due to the Keplerian motion of the disk, the maximum temperature possible in a non-magnetic CV (solid line) is lower by a factor of 2 than in a magnetic CV of the same mass (dashed line, Figure 4). We can use the cooling flow model (Mushotzky & Szymkowiak 1988) to fit the CV spectra, because it is unnecessary to know the post-shock structure to calculate the output spectrum. "Assuming that the same mass flow rate pertains throughout the cooling flow, the emission measure for each temperature is determined by the time it takes for the matter to radiate away sufficient energy to cool down to the next temperature shell. The differential emission measure is thus proportional to the reciprocal of the bolometric luminosity at that temperature."

Figure 5 The Curved Continuum of V1223 Sgr



## 3. "Photoionized" CVs

In Mukai et al. (2003), an underlying power-law continuum was assumed. Such a continuum, however, is unphysical in CVs, and does not fit the data, as clearly demonstrated in the simultaneous Chandra HETG/RXTE PCA observation of V1223 Sgr (Figure 5, where the unfolded data are shown with "best fit" power law, with or without allowing the normalization to be different between the two in-

#### 5. Test Case: V1223 Sgr

I obtain a good fit to the HETG and PCA data on V1223 Sgr above 1 keV by applying pwab and reflect to mkcflow and adding a 6.4 keV Gaussian (Figure 8). The spectral curvature and most of the emission lines are fit well. Below 1 keV, the absorber is more transparent because it is ionized (as evidenced by the OVII edge). The observed Helike lines of O VII, Ne IX, and Mg XI are much stronger than this model: they are indeed likely to be produced via photoionization. Addition of lower temperature plasma does not work because the hypothetical component necessary to explain the Ne IX triplet will also produce additional OVIII lines. Thus, I confirm that photoionization plays a role in V1223 Sgr, creating ionized absorbers and He-like emission lines. Underneath, there is a cooling flow CV modified by the complex absorber.

Figure 7 Resulting Power Law-Like Absorption



Figure 3 The Post-Shock Region



Figure 4 Shock Temperature of non-magnetic CVs



struments). Also in Mukai et al. (2003), all emission lines (except the 6.4 keV fluorescent Fe K $\alpha$ ) were assumed to be due to photoionization: this requires a degree of ionization unlikely in CVs.



4. Complex Absorber Model

The hard continuum in "photoionized" CVs are in fact the cooling flow spectrum modified by a complex absorber. The main absorbers in magnetic CVs are located in the immediate pre-shock flow (see Figure 6), which Done & Magdziarz (1998) approximated as a power law distribution in  $N_H$ . This (pwab)

#### References

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Done, C. & Magdziarz, P. 1998, MNRAS, 298, 737 Mukai, K., Kinkhabwala, A., Peterson, J.R., Kahn, S.M. & Paerels, F. 2003, ApJLett, 486, L77. Mushotzky, R.F. & Szymkowiak, A.E. 1988, in

Channel Energy [keV]

10.0





