

Centaurus A: Interaction of a Radio Source with its Environment

Paul Nulsen

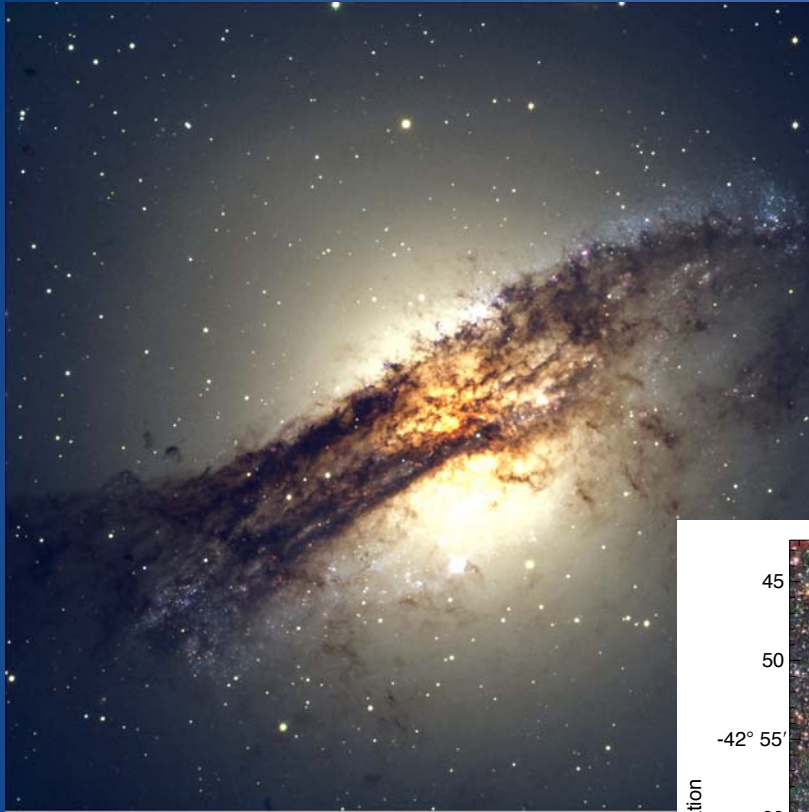
(Harvard-Smithsonian Center for Astrophysics)

Ralph Kraft, David Stark,
JH Croston, MJ Hardcastle,
M Birkinshaw, DM Worrall,
GR Sivakoff, A Jordán, NJ Brassington,
DA Evans, WR Forman, M Gilfanov,
JL Goodger, WE Harris, C Jones, AM Juett,
SS Murray, S Raychaudhury, CL Sarazin,
R Voss, KA Woodley

Centaurus A

Near – distance ≈ 3.7 Mpc (average of 5 values discussed in Ferrarese et al. 2007).

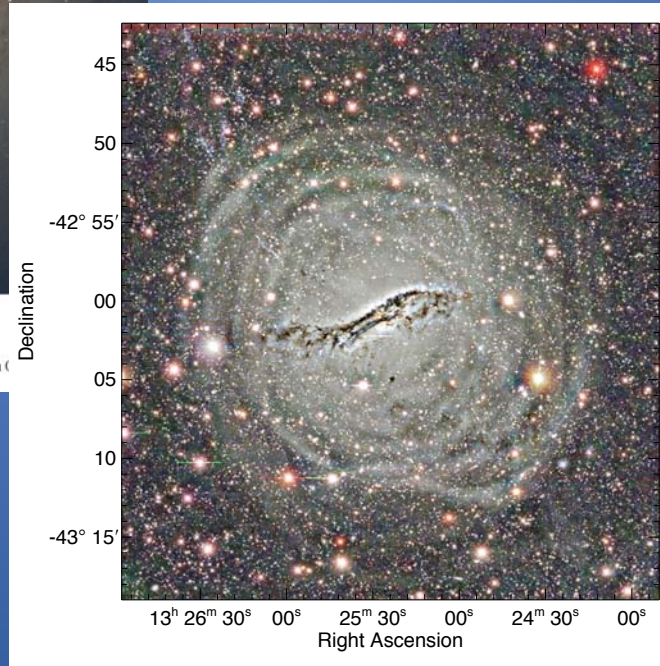
Elliptical galaxy with clear signs of a merger - dust lanes and shells



Centaurus A Radio Galaxy (VLT KUEYEN + FORS2)

ESO PR Photo 05b/00 (8 February 2000)

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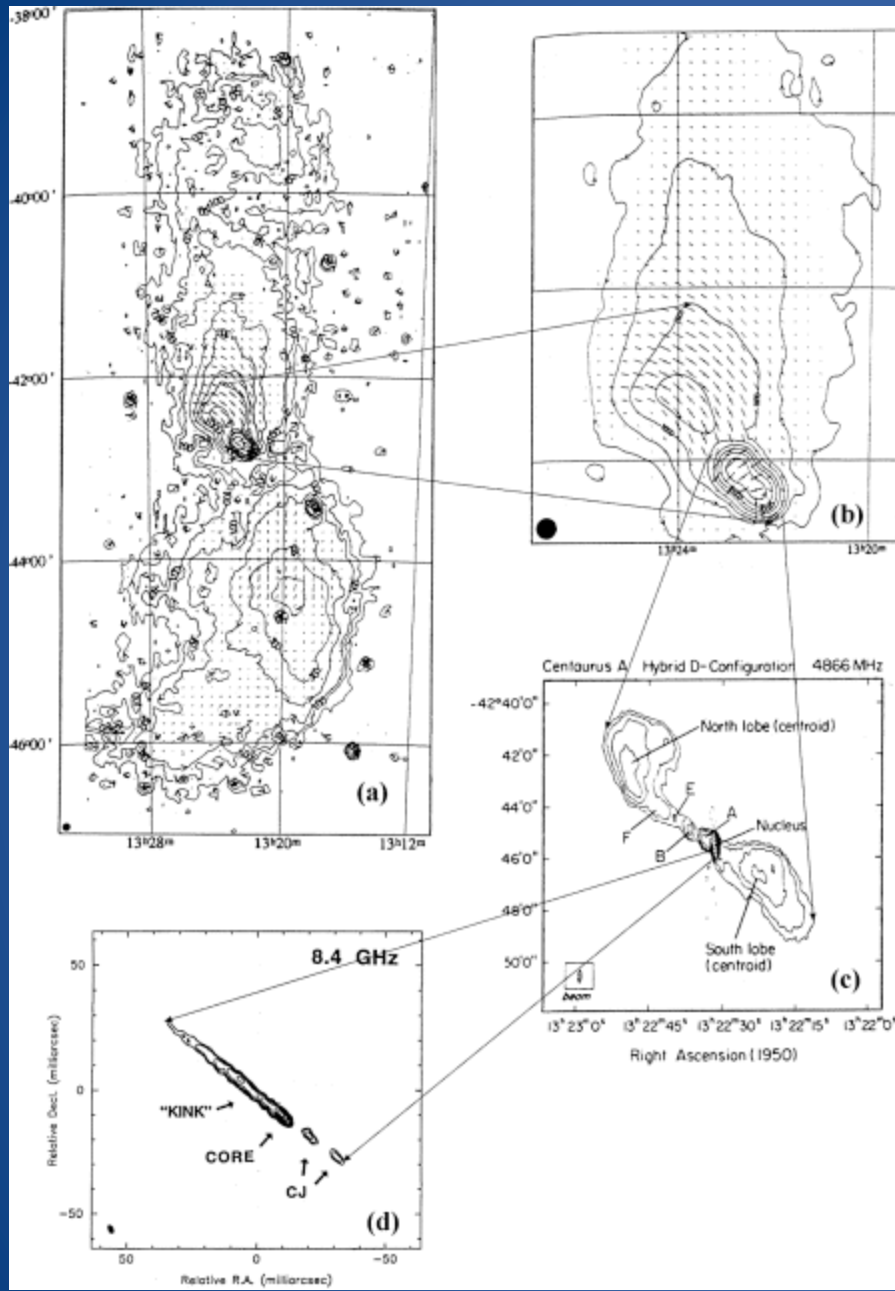
Review by Israel (1998)

Peng et al. 2002

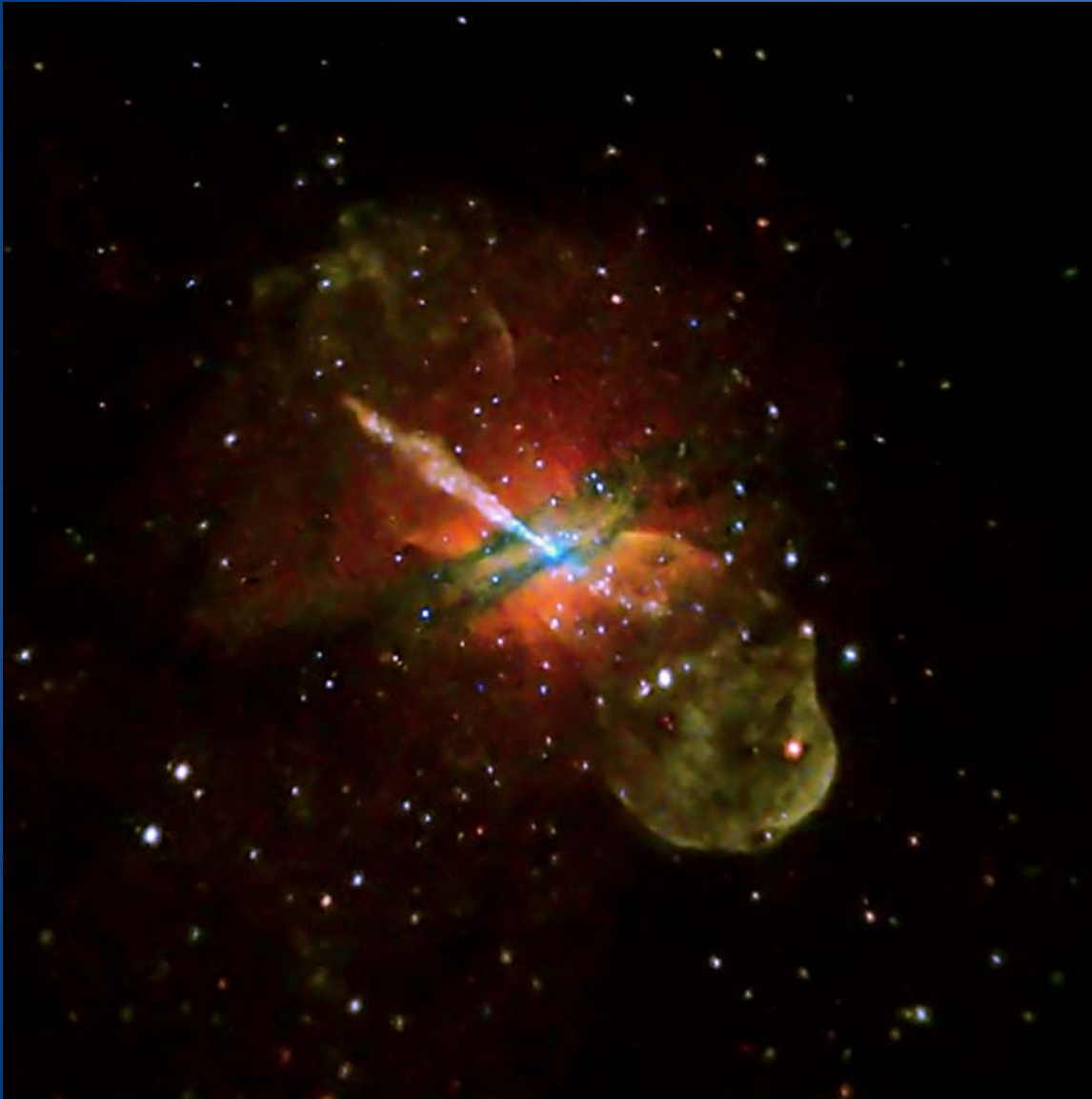
Centaurus A

Nearest extragalactic radio source
- active galaxy

Morganti et al. (1999)



Chandra Image of Centaurus A



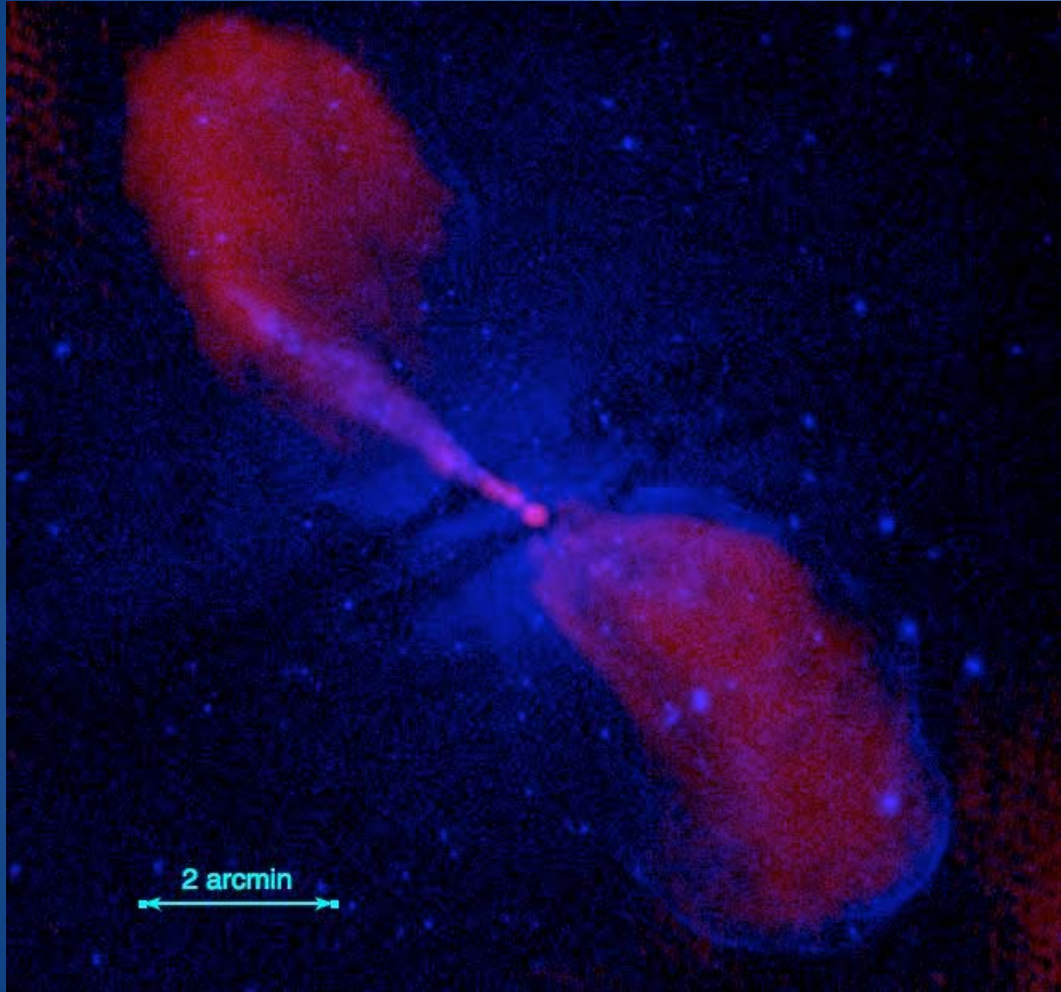
Dust lane

Hot gas

Southwest radio lobe, surrounded
by shock – to ~ 5.5 arcmin (6 kpc)

Jet to northeast

Southwest Radio Lobe

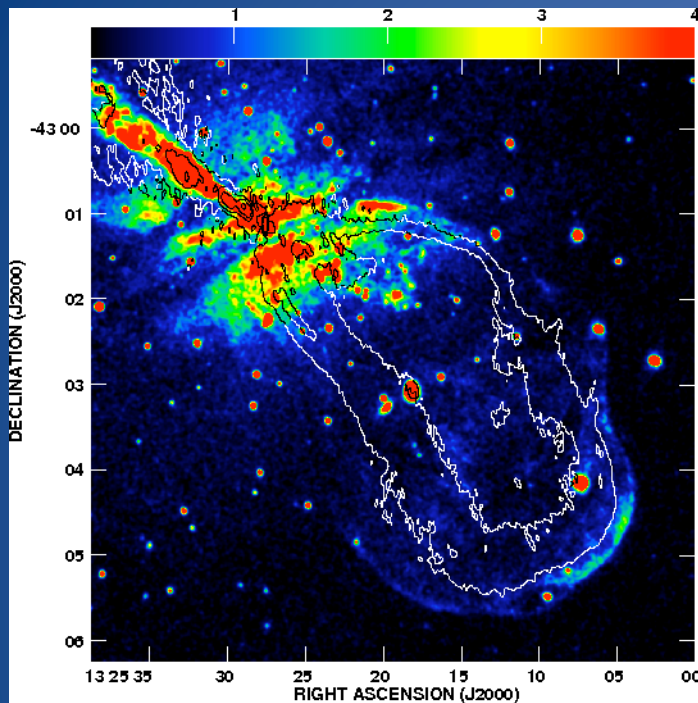


0.3 – 1.5 keV X-ray emission in blue and 5 GHz radio emission in red.

Lobe emission was originally interpreted as shocked ISM (Kraft et al 2003; 2007)

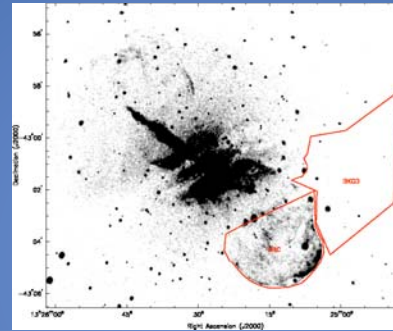
Apparent density jump exceeded 4, maximum compression for hydrodynamic shock with ratio of specific heats $\Gamma = 5/3$.

Synchrotron Shock Model



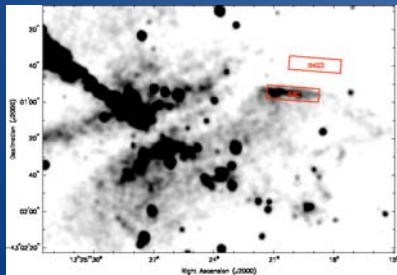
Croston et al (2009):

X-ray spectral fits to regions around the SW lobe strongly prefer power-law to thermal models.

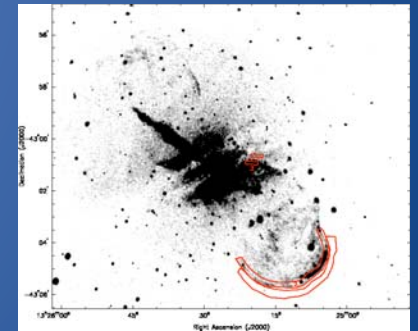


Power law fit: $\chi^2 = 931/744$ dof vs worse than 1010/743 for thermal models (adding thermal component gave insignificant improvement). No line emission.

Power law fit: $\chi^2 = 333/309$ dof vs worse than 344/308 for thermal models



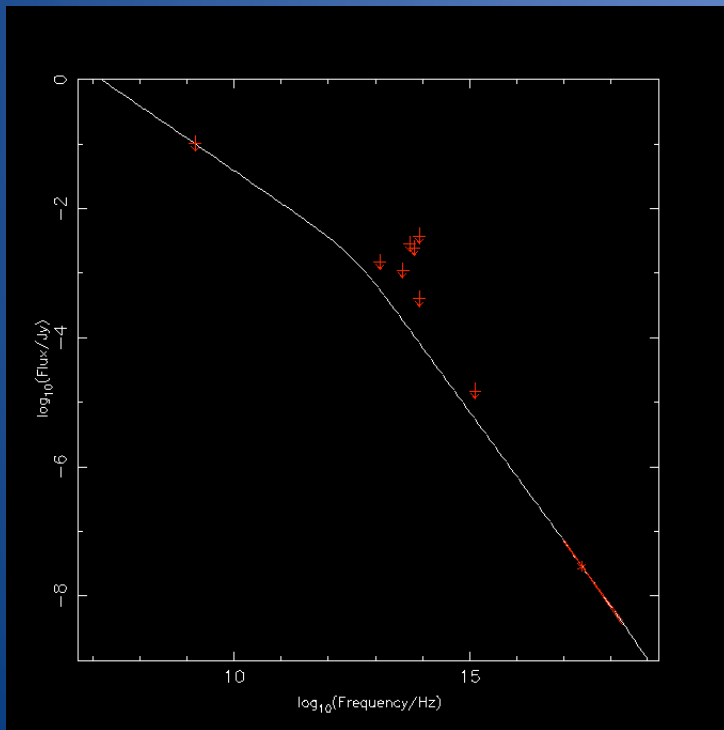
Thermal model is preferred for region closer to the nucleus: $\chi^2 = 137/97$, $kT \approx 1$ keV, vs 445/98 for power law model. Emission lines prominent.



Synchrotron Shock Model

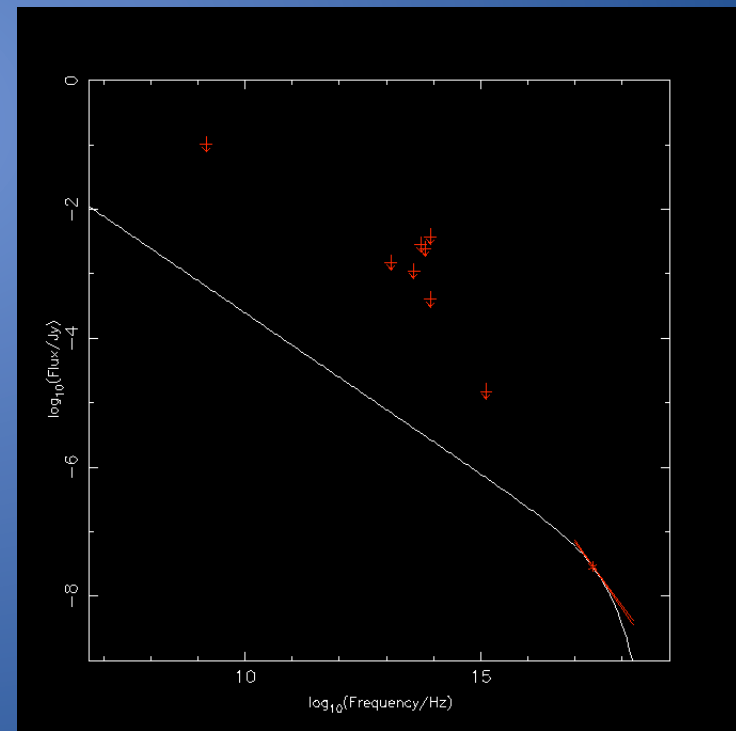
Expect synchrotron emission across a wide range of frequencies. Reasonable models for the electron energy distribution can produce the X-ray emission without exceeding upper limits at other wavelengths. (IC models inconsistent with lack of radio emission.)

$$\alpha = 0.5 \ (p = 2), \ \gamma_{\min} = 10, \\ \gamma_{\text{break}} = 4.4 \times 10^4, \ \gamma_{\max} = 5 \times 10^9.$$



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$$\alpha = 0.5 \ (p = 2), \ \gamma_{\min} = 10, \\ \text{no break}, \ \gamma_{\max} = 3 \times 10^8.$$



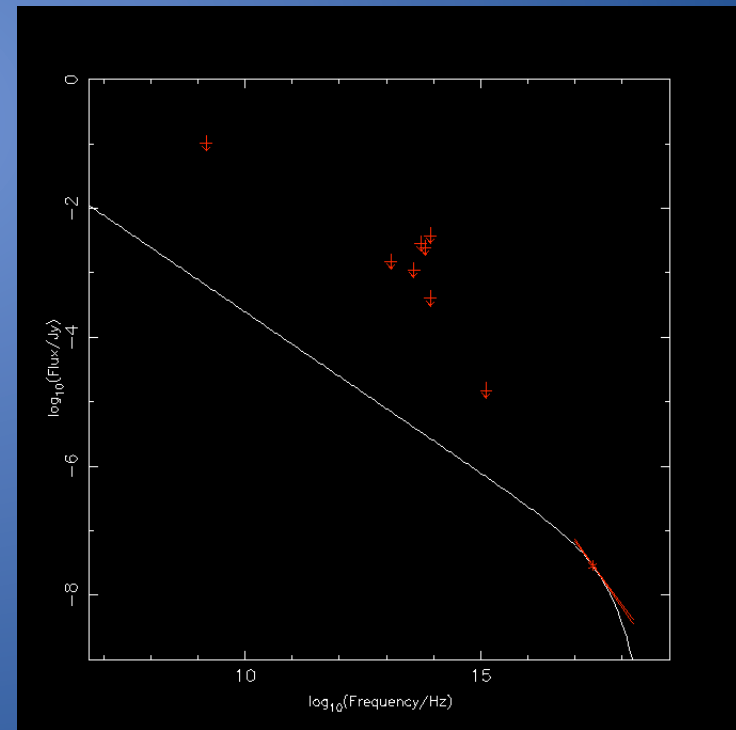
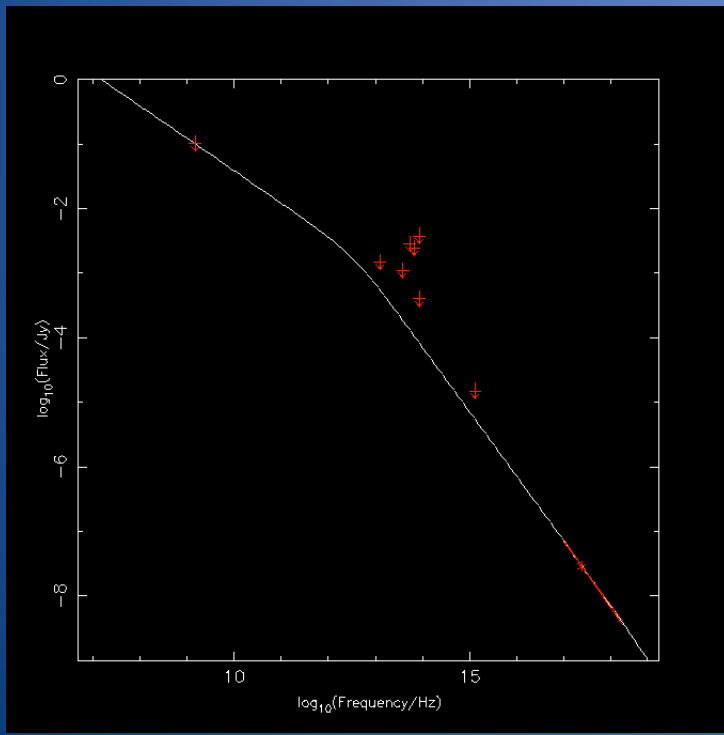
Chandra's 1st Decade of Discovery

Synchrotron Shock Model

Models for diffusive shock acceleration, where γ_{\max} is determined by competition between acceleration timescale and radiative losses (eg Reynolds 1996), imply $\gamma_{\max} \sim 10^8$ at the tip of SW lobe (where the shock is strongest). Variation of shock speed may then account for the lack of non-thermal X-ray emission from the shock nearer the AGN.

(Upper limit) $\alpha = 0.5$ ($p = 2$), $\gamma_{\min} = 10$,
 $\gamma_{\text{break}} = 4.4 \times 10^4$, $\gamma_{\max} = 5 \times 10^9$.

$\alpha = 0.5$ ($p = 2$), $\gamma_{\min} = 10$,
no break, $\gamma_{\max} = 3 \times 10^8$.



Jet Power

Behind the thermal shock in north:

Proton density $n_p = 0.033 \text{ cm}^{-3}$ and $kT = 0.95 \text{ keV} \Rightarrow \text{pressure} = 1.1 \times 10^{-10} \text{ cgs}$.

High sound speed \Rightarrow pressure nearly uniform in lobe, so same pressure drives SW shock .

Outside shock, $n_e \approx 0.001 \text{ cm}^{-3}$, $kT = 0.35 \text{ keV}$ (Kraft et al 2003), so pressure jump at shock $\approx 87 \Rightarrow$ shock Mach number of ≈ 8.4 , or $v_{\text{shock}} \approx 2600 \text{ km s}^{-1}$.

Thermal emission from the shocked gas is negligible compared to the non-thermal emission observed.

Integrating over energy distribution of non-thermal electrons and including equipartition magnetic field (no protons) gives their contribution to the pressure as no more than $\approx 2 \times 10^{-12} \text{ cgs}$ – negligible compared to the thermal pressure.

Shock age, $2a/v_{\text{shock}} \approx 2 \times 10^6 \text{ y} \Rightarrow \text{average power} \approx 4pV/\text{age} \approx 10^{43} \text{ erg s}^{-1}$.

“Instantaneous” power, $P = p \, dV/dt = pV(1/a \, da/dt + 2/b \, db/dt) = 3pV \, 1/a \, da/dt = 4\pi r b^2 \, da/dt = 2\pi r b^2 v_{\text{shock}} \approx 6.6 \times 10^{42} \text{ erg s}^{-1}$.

NE Jet Flow Model

Like Laing & Bridle (2002):

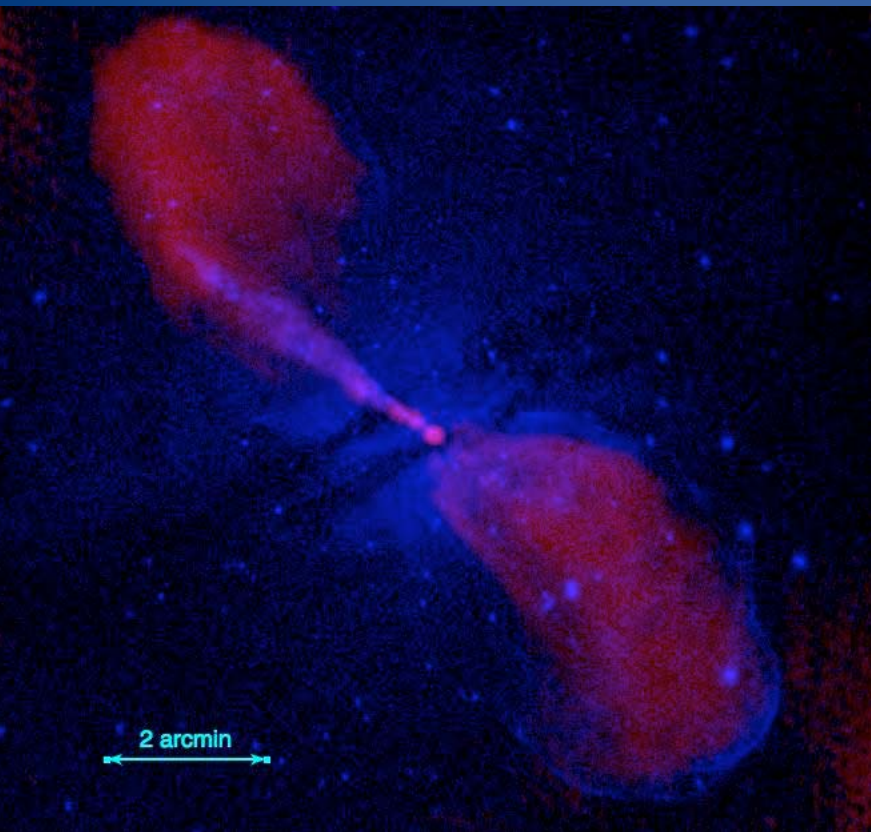
Gas temperature ($kT \approx 0.55$ keV), density profile ($n_e \sim r^{-1.26}$), pressure from Chandra data . **Equate to internal pressure.

Steady, near 1-d flow. Area of cross, A , varies with radius, R .

Proper density of jet rest mass, ρ . Rate of flow through jet $\dot{M} = \gamma\beta c A \rho$

Allowing for entrainment α is rate of mass injection

$$\dot{M} \Big|_1^2 = \int_1^2 \alpha A dR$$



Cen A, 5 GHz (red) and 0.3 – 1.5 keV (blue)

Power through jet assumed constant:

$$P = (\gamma - 1) \dot{M} c^2 + h A c \beta \gamma^2$$

h = enthalpy per unit volume, $h = p + e = \Gamma p / (\Gamma - 1)$, for pressure p . Here $\Gamma = 13/9$.

Momentum flux

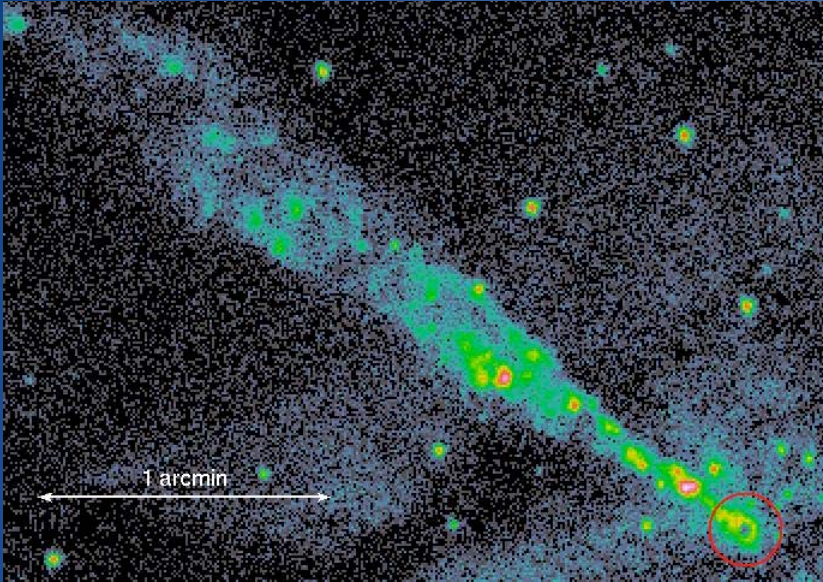
$$\Pi = (P/c + \dot{M} c) \beta$$

affected by buoyancy

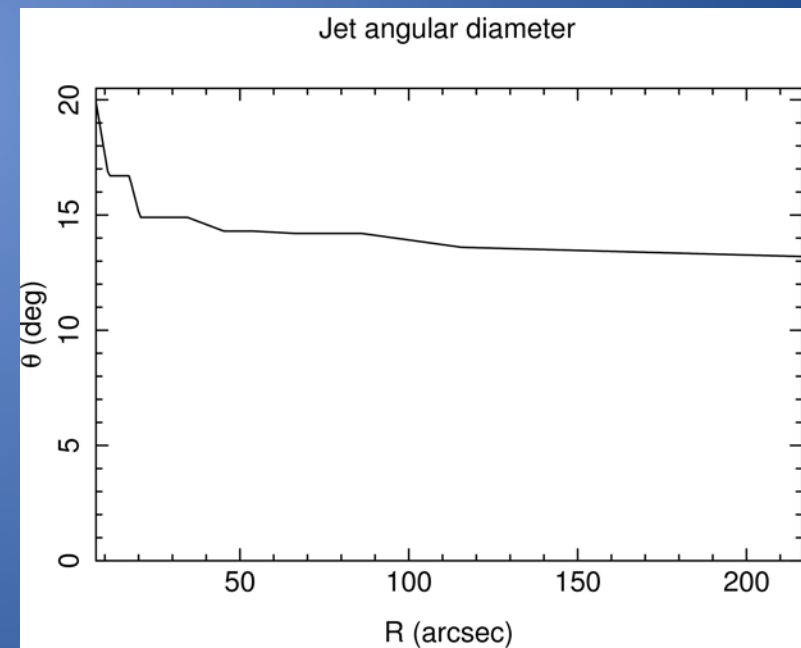
$$\Pi \Big|_1^2 = \int_1^2 \frac{dp}{dR} A dR$$

Jet Width

Knots complicate measurement of jet width



Plot shows angle subtended by jet at the AGN



Flow Parameters

Fiducial values:

Jet power, $P = 6 \times 10^{42} \text{ erg s}^{-1}$ (cf $6.6 \times 10^{42} \text{ erg s}^{-1}$ for power into SW lobe, Croston et al 2009).

Initial speed, $\beta = 0.7$: radio knot proper motions of $\approx 0.5c$ near start of jet (Hardcastle et al 2003).

Mass injection:

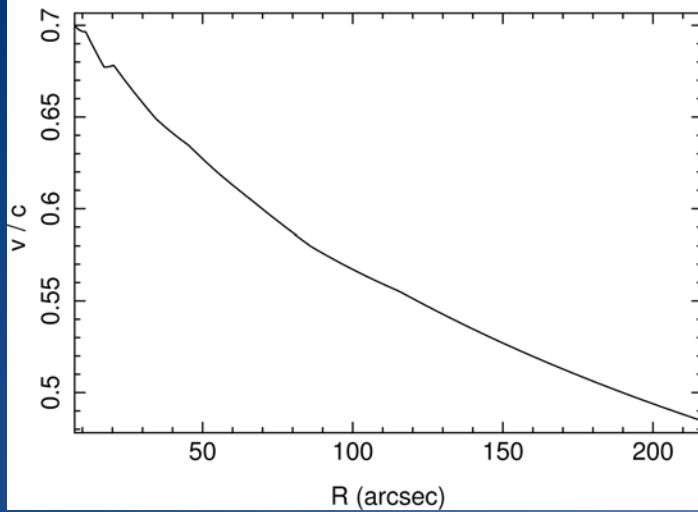
Star density is expressed as a fraction, f , of gravitating mass density (determined from hydrostatic equilibrium). Expect $f \approx 1$ at $R = 100 \text{ arcsec}$ (1.8 kpc; also consistent with photometry). Then $\alpha = f \rho_{\text{grav}} / \tau$, with $\tau = 10^{12} \text{ yr}$ (Faber & Gallagher 1976).

Allow for other entrainment by varying f .

In fact, solution is over-determined, so we adjust f to make change in momentum flux match the buoyant force.

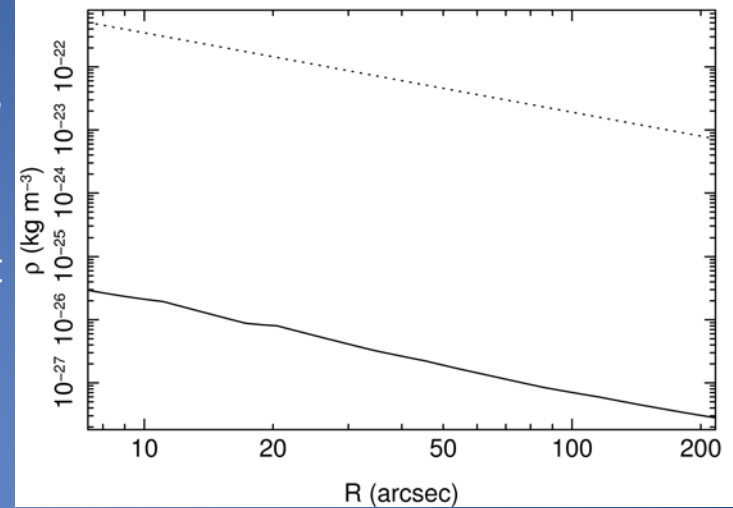
Fiducial Model

Jet beta

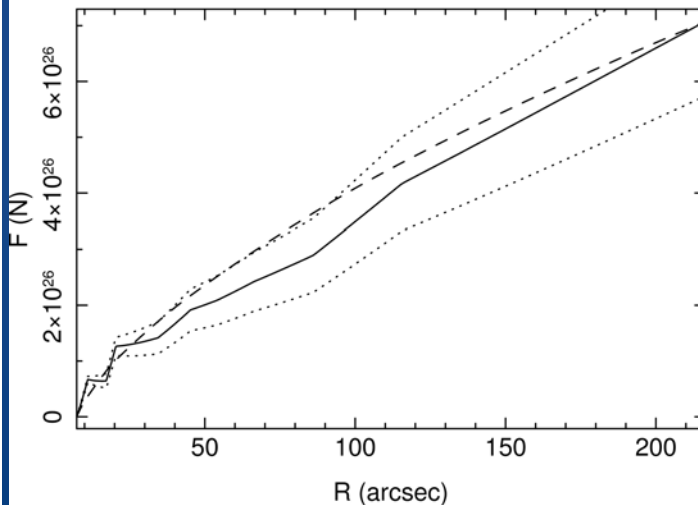


Non-dissipative flow would obey Bernoulli's theorem: pressure decrease => speed increase – inconsistent with data.

Jet and gas density



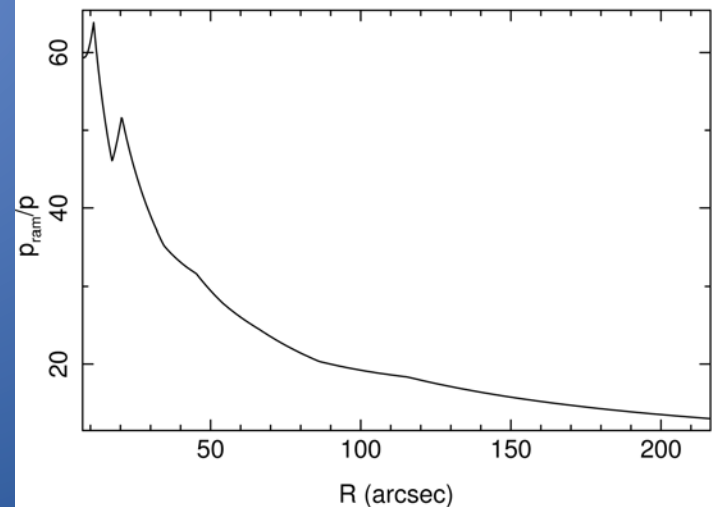
Change in momentum flux and buoyant force



Need $f = 0.60$ to make $\Delta\Pi = \text{buoyant force}$ -

ie, stellar mass loss within the jet can account for all dissipation.

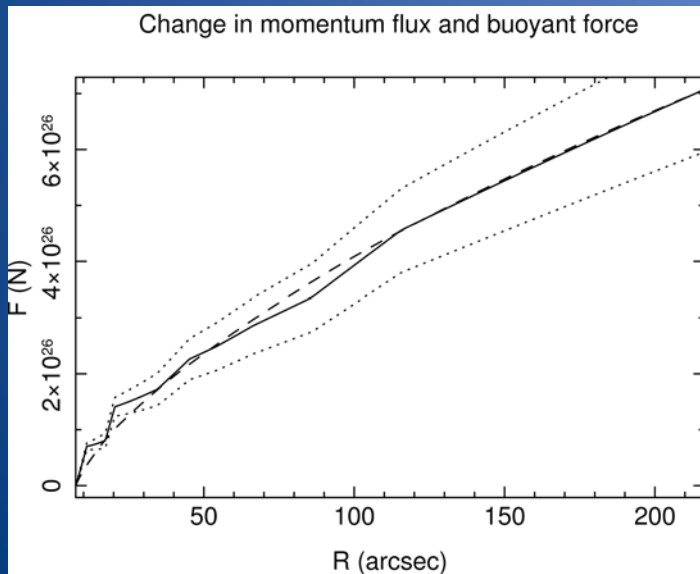
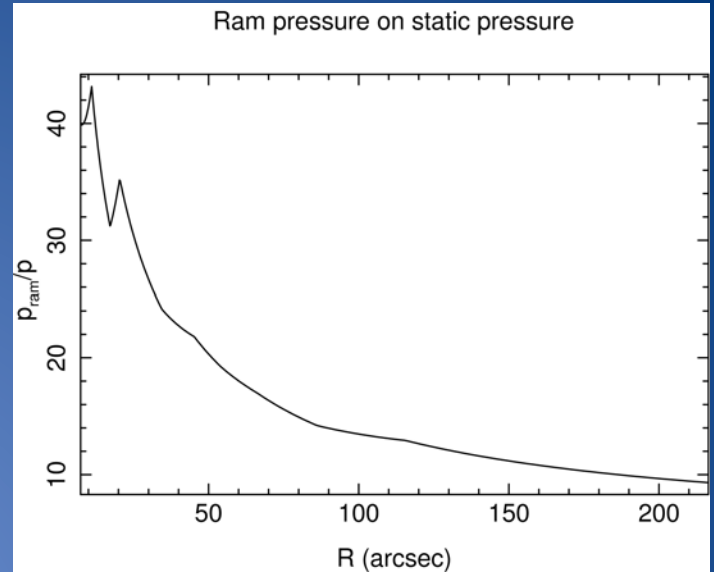
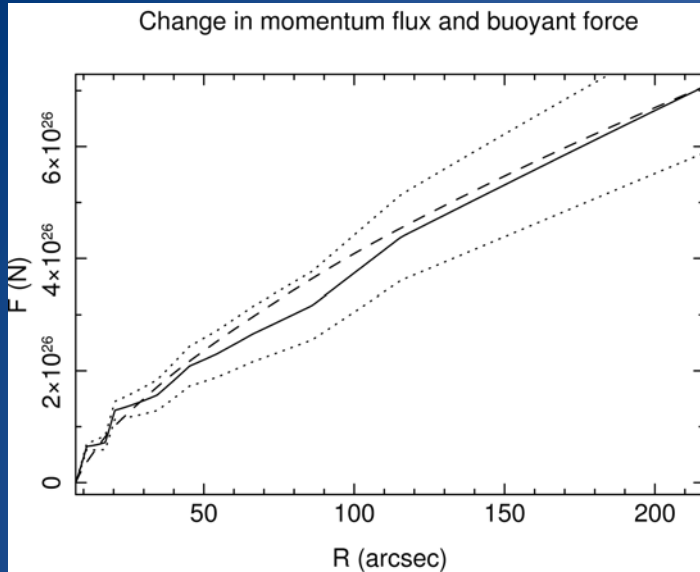
Ram pressure on static pressure



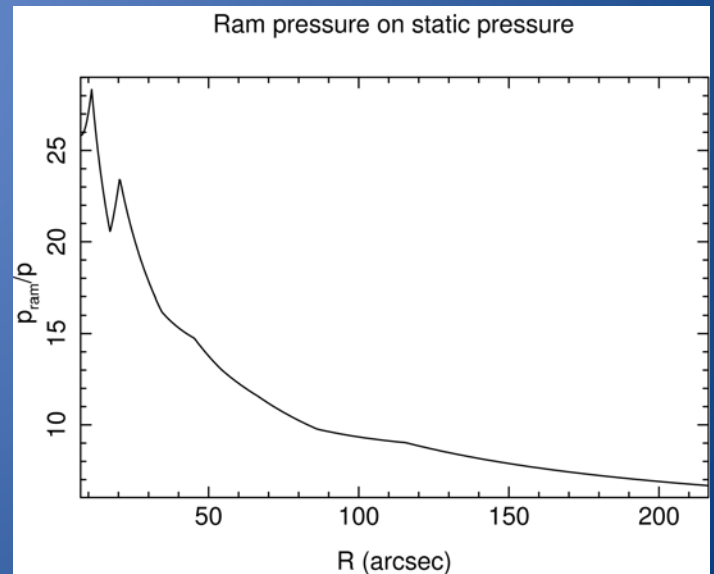
Variations

Initial $\beta = 0.85$
need $f = 0.42$

(for $\beta = 0.5$, $f = 0.93$)



$P = 3 \times 10^{42} \text{ erg s}^{-1}$
need $f = 0.61$

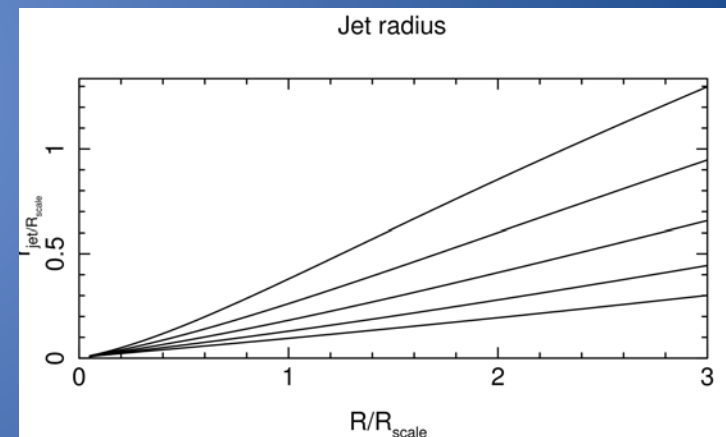
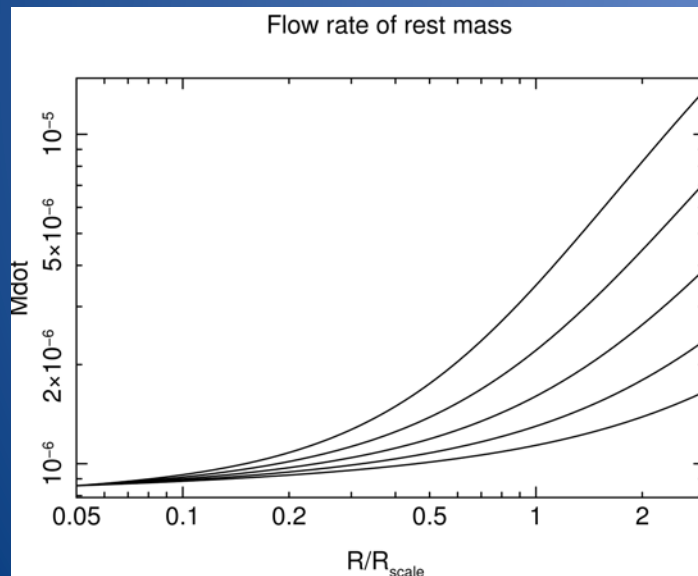


Effect of Environment

If dissipation is governed by stellar mass injection, the jet is unstable:
larger cross section => more stars in jet, hence more dissipation => more broadening

Model:
$$\frac{d\dot{M}}{dR} = \alpha A; \quad \frac{d\Pi}{dR} = -A \frac{dp}{dR}; \quad P = \text{constant}$$

Same power, initial speed, mass injection (α) as fiducial model, but bounding pressure is scaled down by a factor up to 2:



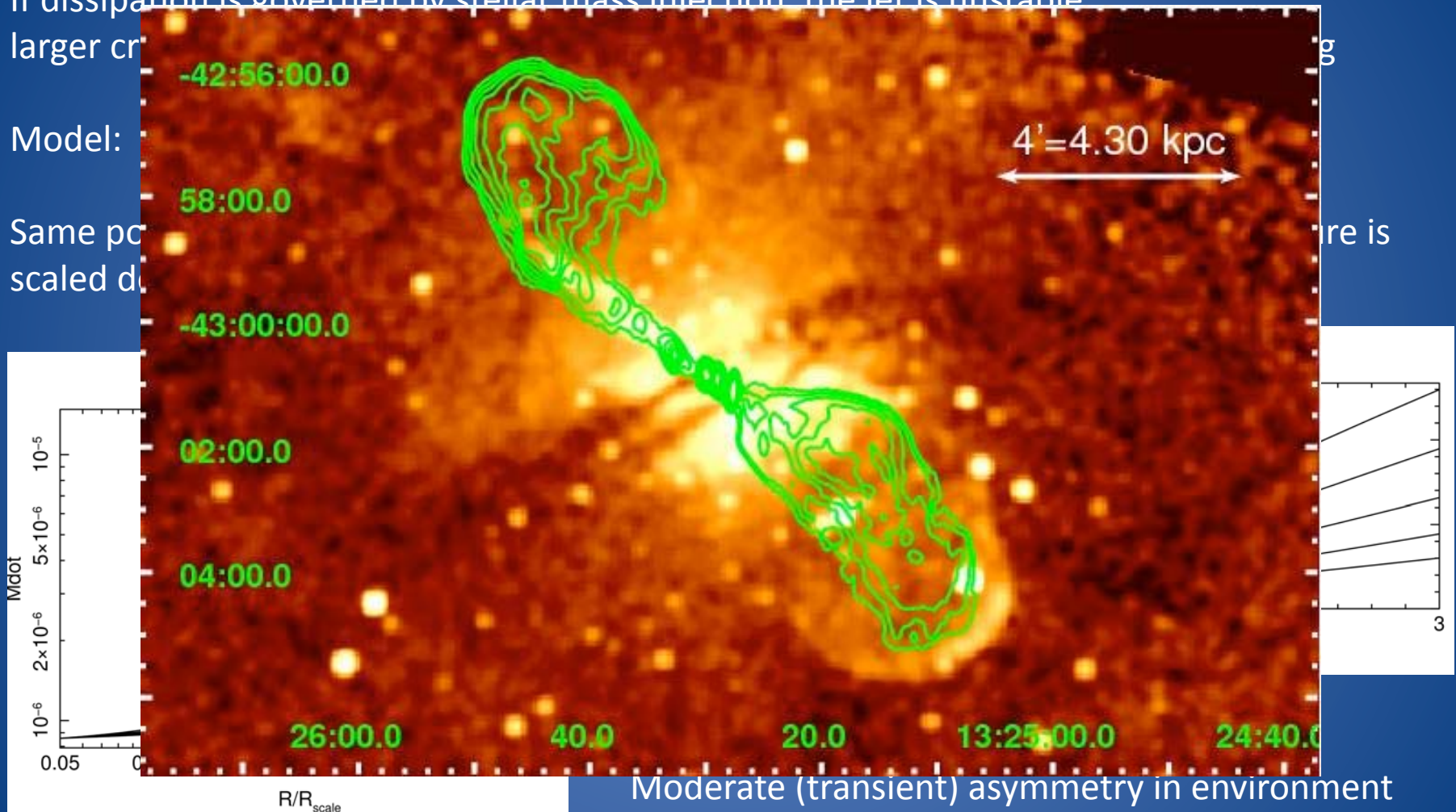
Moderate (transient) asymmetry in environment may make the difference between a jet and a lobe.

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larger cr

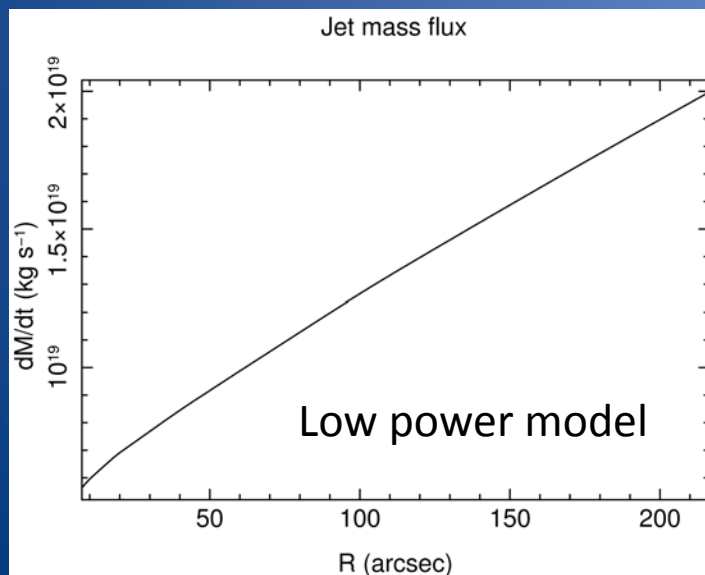
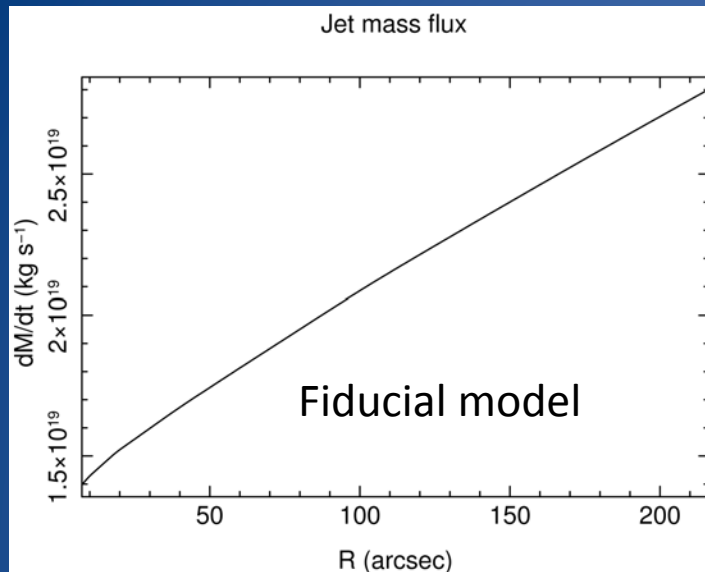
Model:

Same po
scaled d



Moderate (transient) asymmetry in environment
may make the difference between a jet and a lobe.

Stellar Mass Loss



Total mass injection $\approx 2 \times 10^{-4} M_{\odot} \text{ yr}^{-1}$, comparable to mass loss rate of ~ 100 AGB stars.

Comparable to the number of knots in the jet (eg Hardcastle et al 2007), suggesting that each knot may be caused by a single star.

Ram pressure stops wind where $p_{\text{ram}} \approx \rho_{\text{wind}} v_{\text{wind}}^2$, ie for a star losing $10^{-6} M_{\odot} \text{ yr}^{-1}$, with $v_{\text{wind}} = 10 \text{ km s}^{-1}$, at $20''$ from the AGN, only $\approx 0.009 \text{ pc}$ from the star. (cf. $\sim 3 \text{ pc}$, Tingay & Lenc 2009).

A wind of $\approx 10^{-6} M_{\odot} \text{ yr}^{-1}$ must intercept $\approx 1\%$ of jet (mass flux $\approx 10^{-4} M_{\odot} \text{ yr}^{-1}$) to be accelerated to jet speed

\Rightarrow long trails behind AGB stars where wind gas is accelerated and mixed with jet gas

\Rightarrow supersonic, turbulent wake?

Conclusions

- X-ray emission from shock surrounding SW radio lobe is predominantly synchrotron
- The shock speed $\approx 2600 \text{ km s}^{-1}$ and the jet power $\approx 6.6 \times 10^{42} \text{ erg s}^{-1}$
- Stellar mass loss may be the primary source of mass entrained by the Cen A jet
- Modest asymmetry in the environment can make the difference between a jet and a lobe
- Many jet knots may be due to individual AGB stars within the jet